

# **Fiber-Base duality, Global Symmetry Enhancement and Gopakumar-Vafa invariant**

Futoshi Yagi (Technion)

Based on

arXiv: 1411.2450: V. Mitev, E.Pomoni, M.Taki, FY

Work in progress: H.Hayashi, S-S.Kim, K.Lee, M.Taki, FY

**5D  $N=1$  SUSY SU(2) gauge theory  
with  $N_f$  flavor**

# 5D $N=1$ SUSY SU(2) gauge theory with $N_f$ flavor

**5D UV fixed point exists for  $N_f \leq 7$**

'96 Seiberg

# 5D $N=1$ SUSY $SU(2)$ gauge theory with $N_f$ flavor

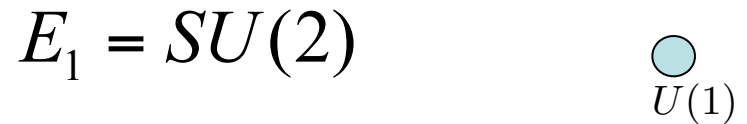
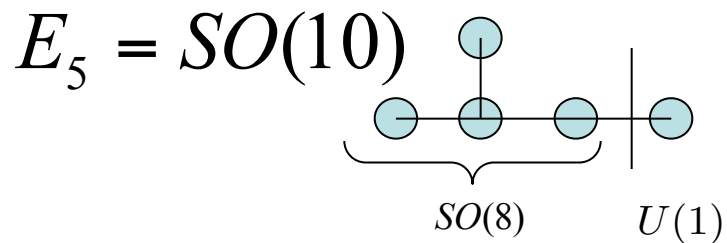
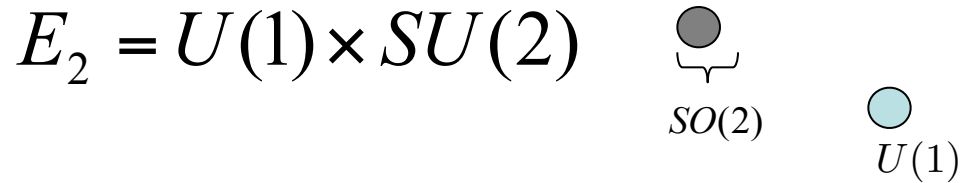
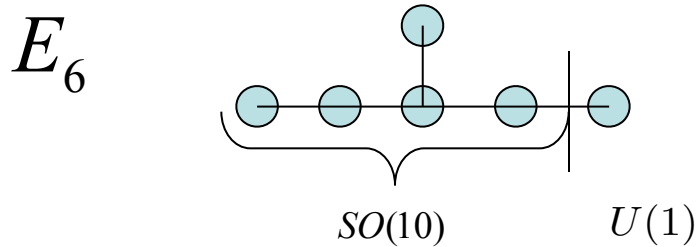
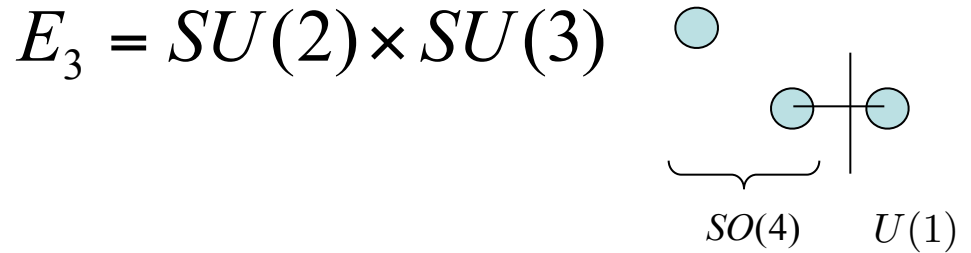
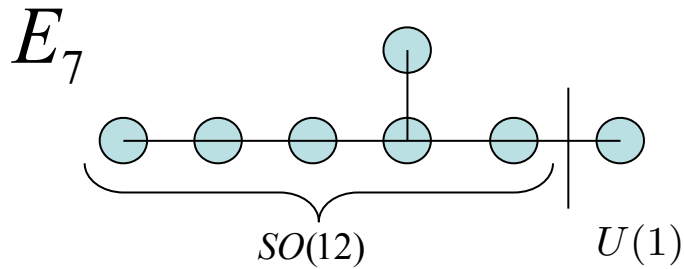
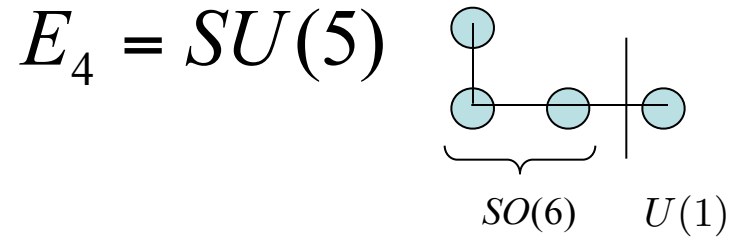
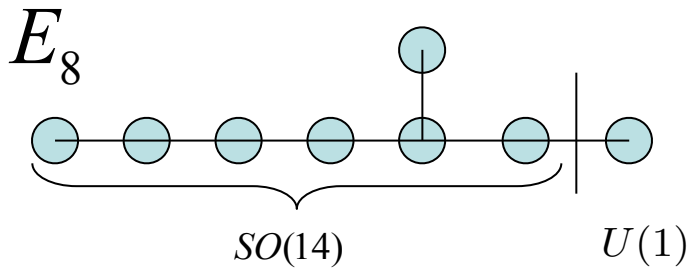
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**Global symmetry enhancement  
at UV fixed point**

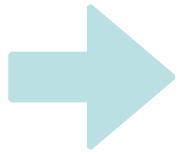
$$SO(2N_f) \times U(1) \subset E_{N_f+1}$$

$N_f$  flavors                      Instanton particle



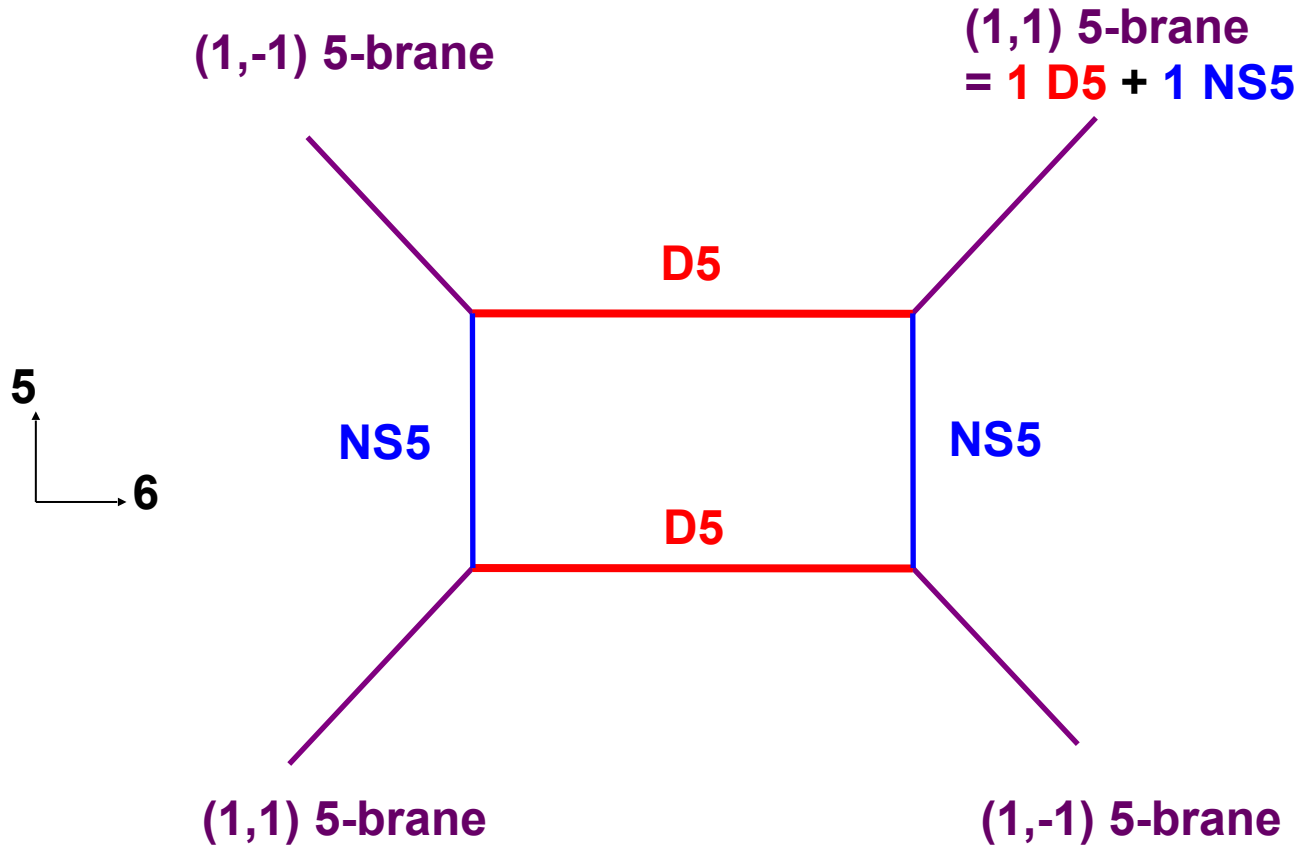
**Can we see global symmetry  
enhancement from brane web?**

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**S-duality**  
**(Fiber-base duality in CY language)**

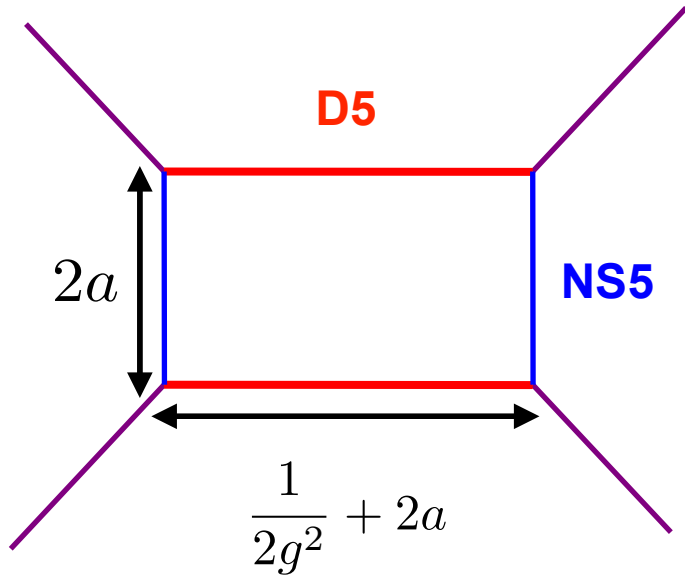
# Brane setup for pure SU(2) SYM



NS5	0	1	2	3	4	5
D5	0	1	2	3	4	6

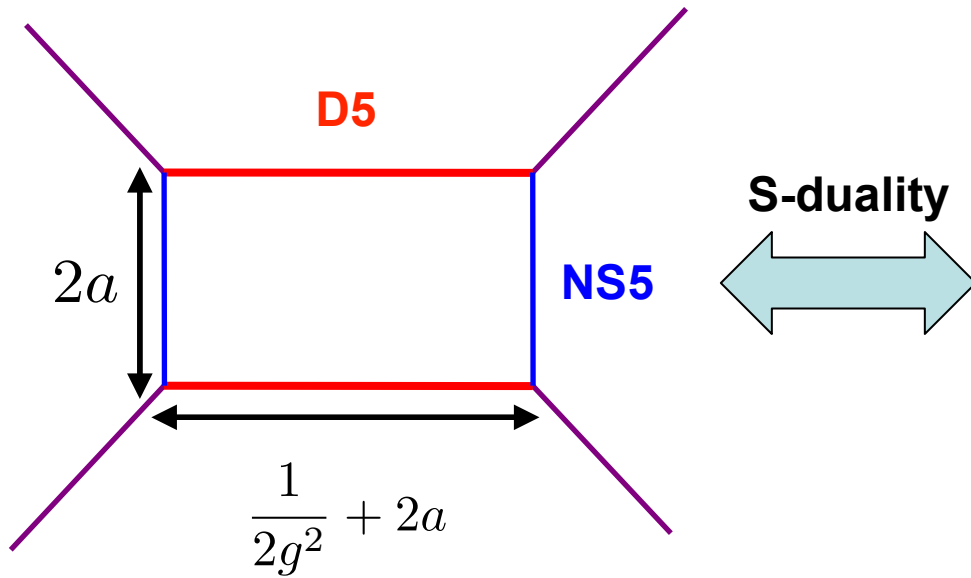


# pure SU(2) SYM

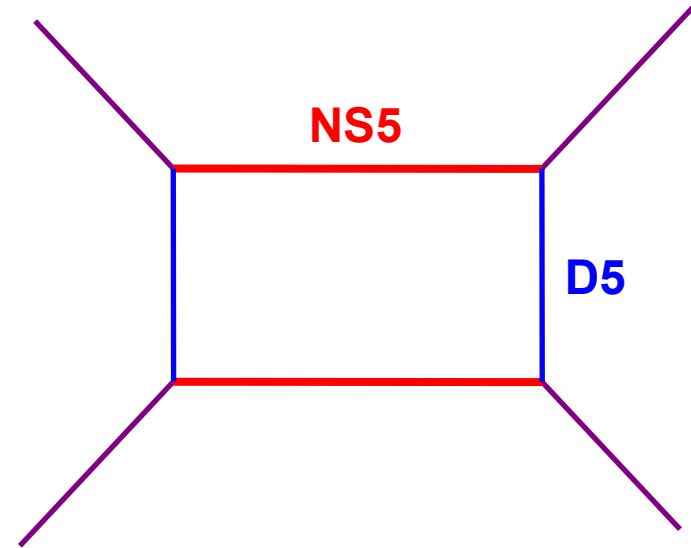



$a$  : Coulomb moduli parameter  
 $g$  : (Bare) gauge coupling

# S-duality for pure SU(2) SYM

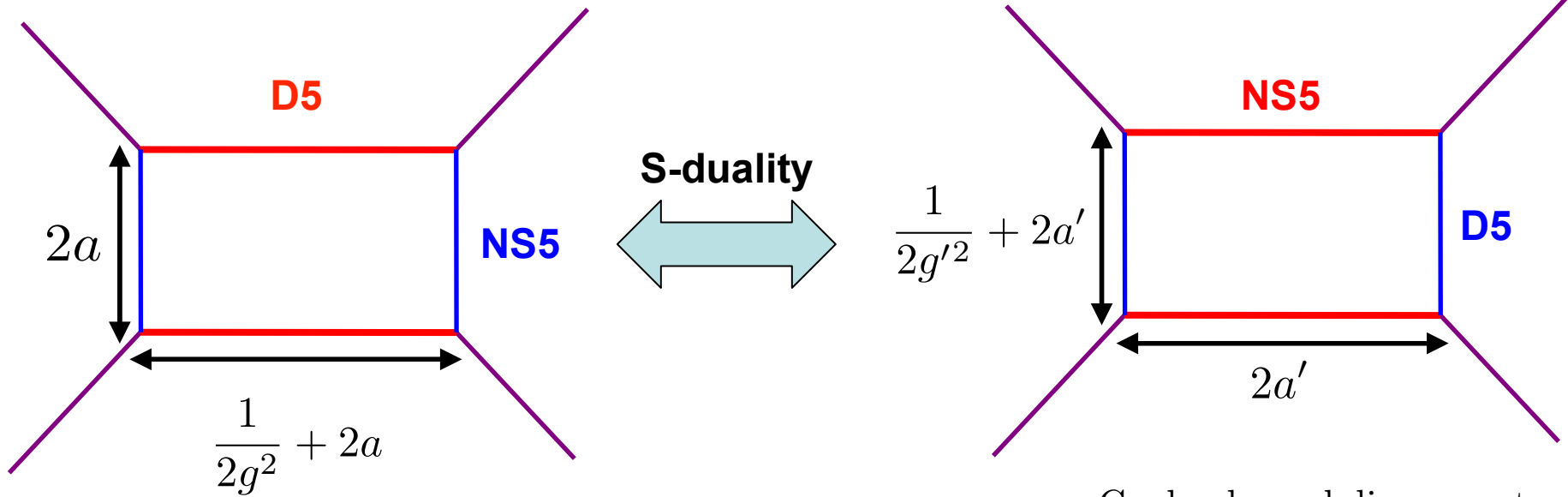


**S-duality**



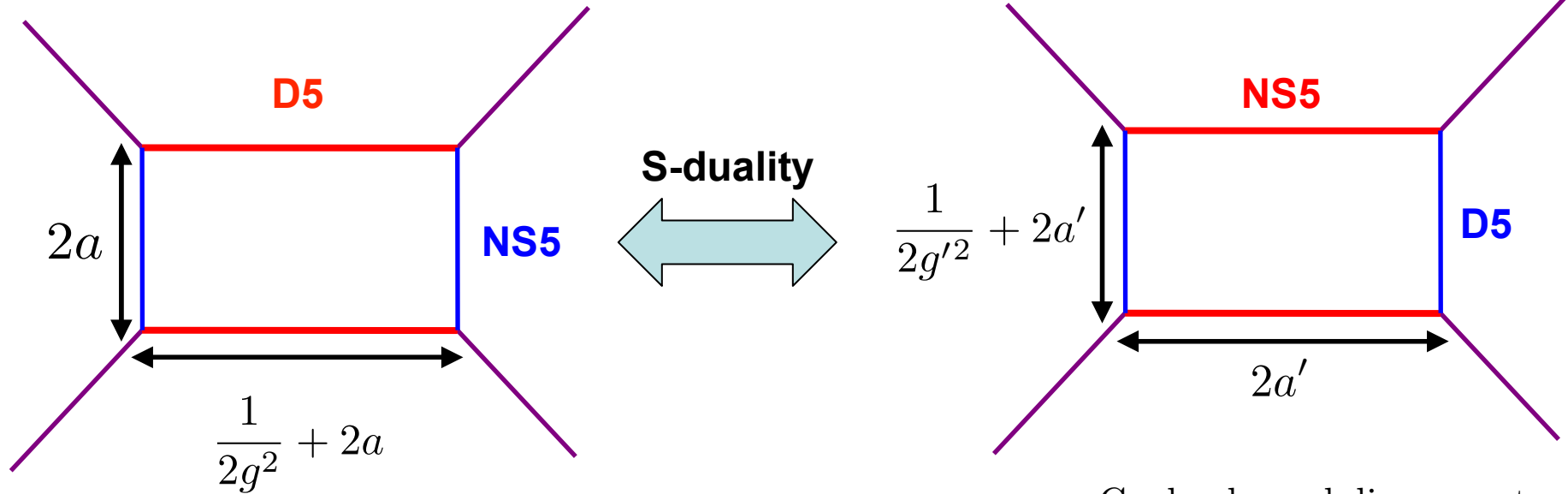
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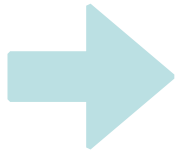


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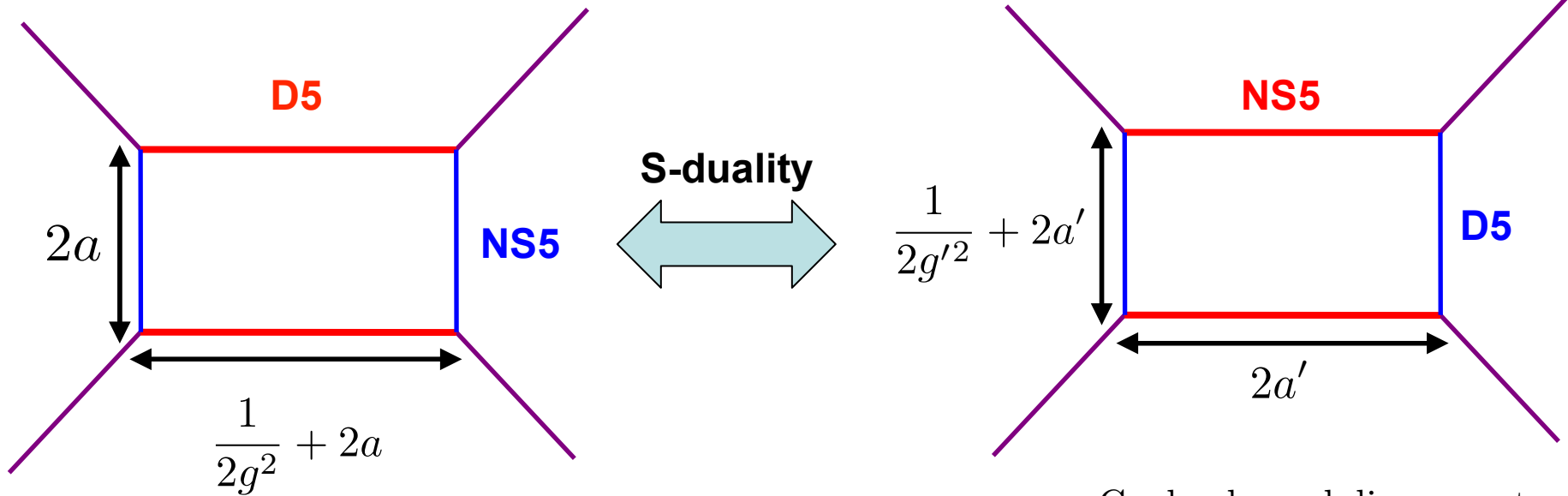
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$$\frac{1}{g'^2} = -\frac{1}{g^2}$$

$$a' = a + \frac{1}{4g^2}$$

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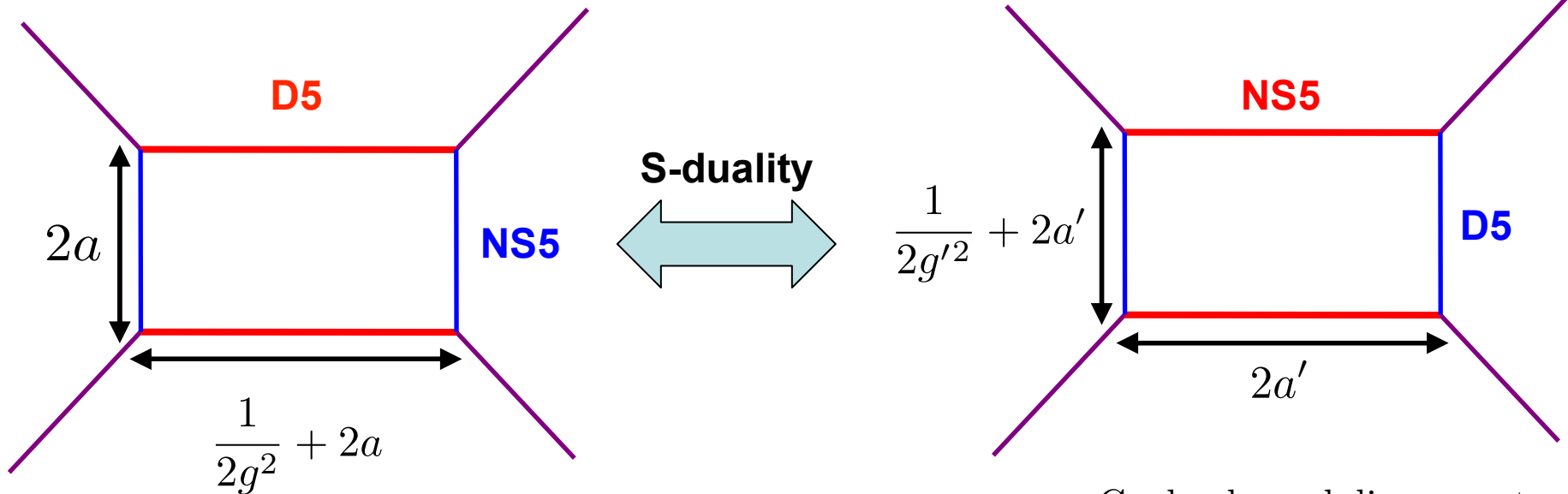
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**Weyl Symmetry for  $E_1 = SU(2)$**

'97 Aharony, Hanany, Kol

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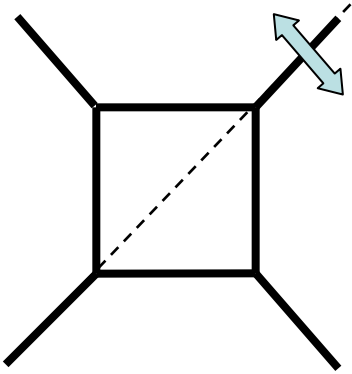
**Weyl Symmetry for  $E_1 = SU(2)$**

'97 Aharony, Hanany, Kol

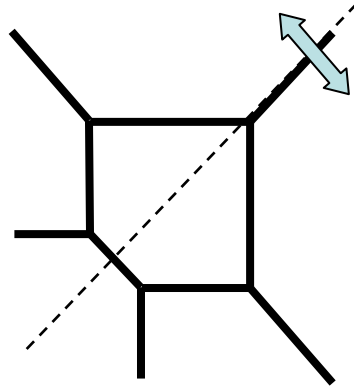
**Coulomb moduli parameter is also transformed!**

# Generalization to higher flavor

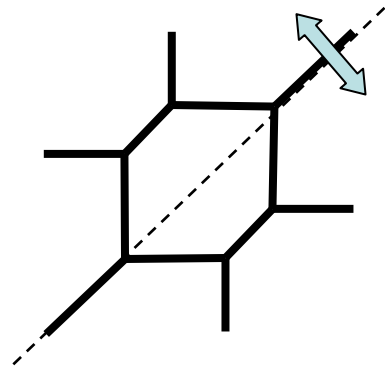
$N_f = 0$



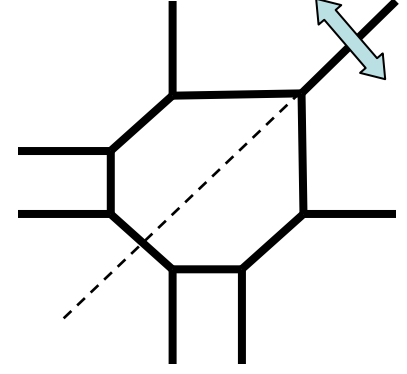
$N_f = 1$



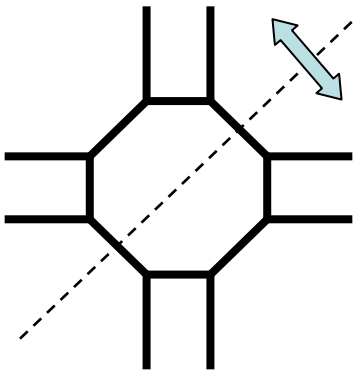
$N_f = 2$



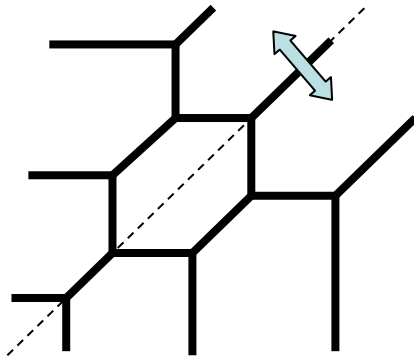
$N_f = 3$



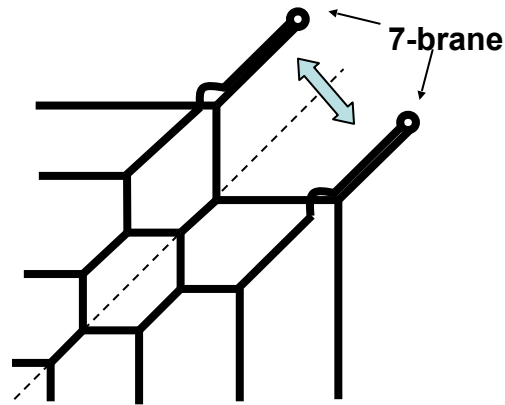
$N_f = 4$



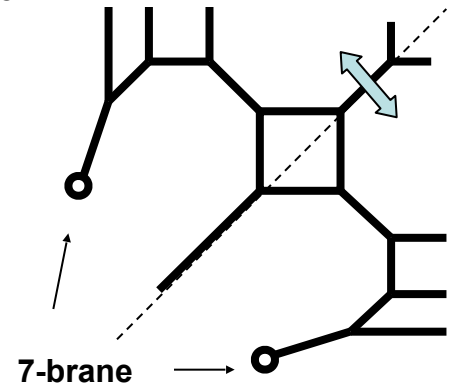
$N_f = 5$



$N_f = 6$



$N_f = 7$



(Weyl symmetry of)  
 **$SO(2N_f) \times U(1)$**

+

Transformation induced from  
**S-duality**



(Weyl symmetry of)

$$\mathbf{SO}(2N_f) \times \mathbf{U}(1)$$

+

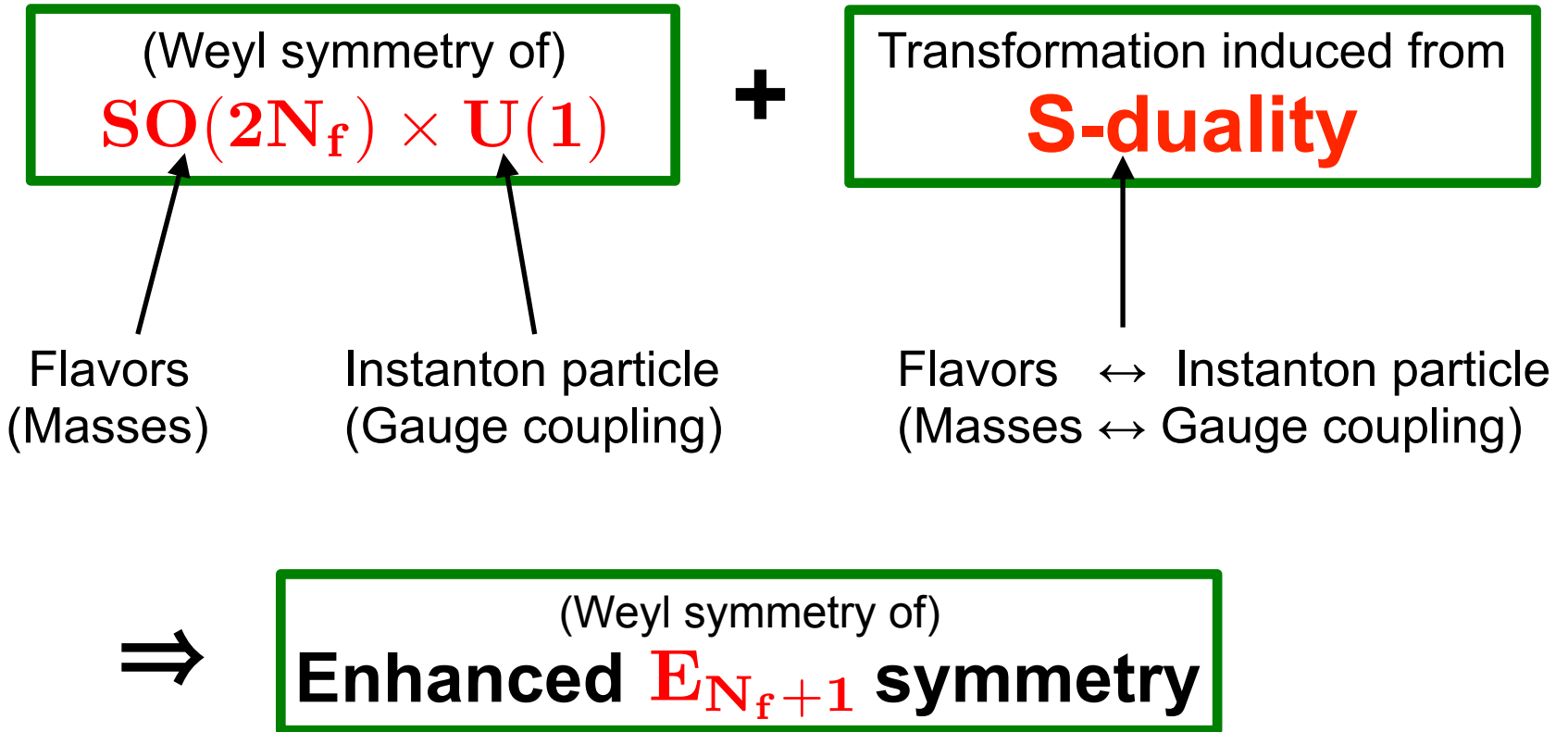
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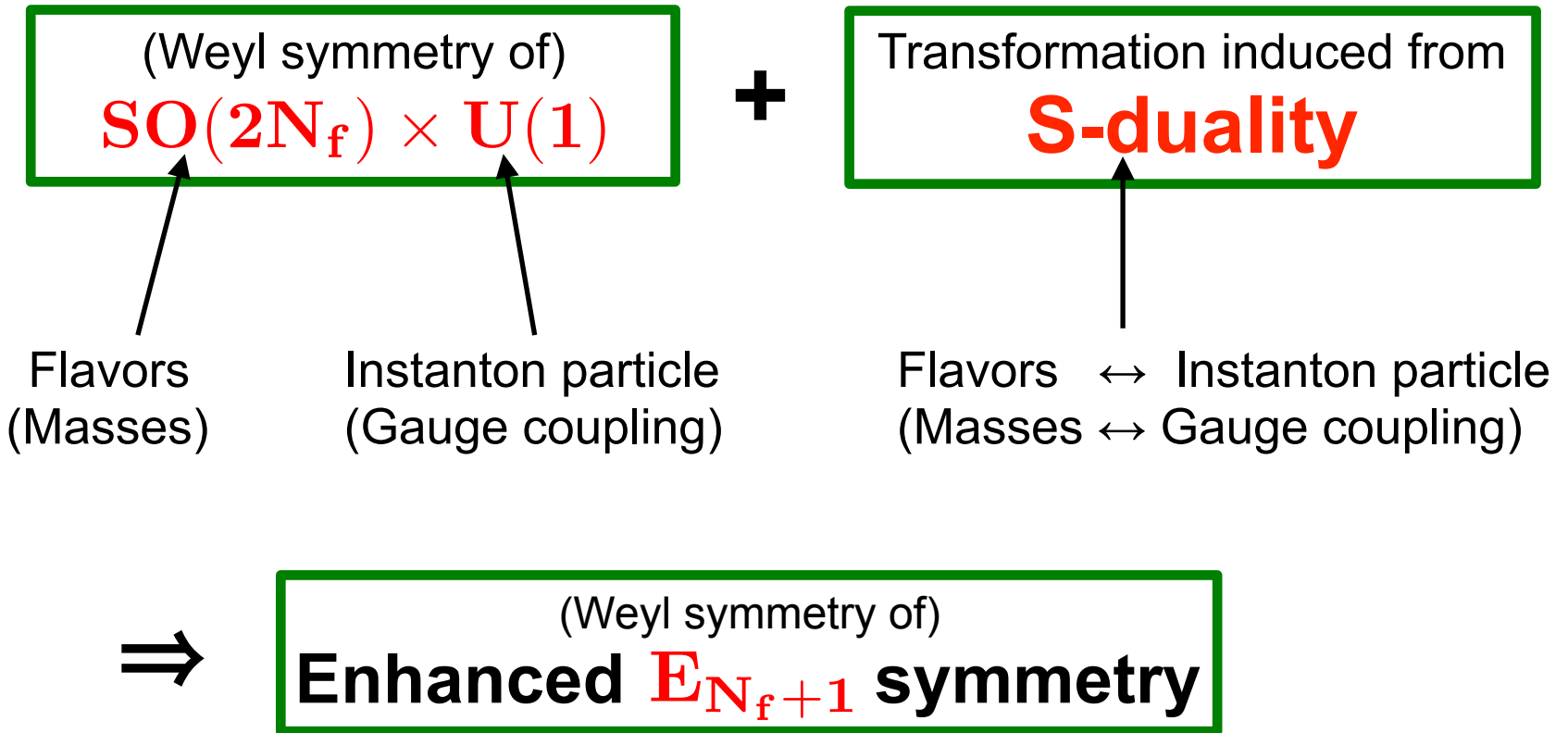
**S-duality**

Flavors  
(Masses)

Instanton particle  
(Gauge coupling)

Flavors ↔ Instanton particle  
(Masses ↔ Gauge coupling)





**Again, Coulomb moduli parameter is also transformed!**

**Can we write Nekrasov partition function  
in manifestly  $E_{Nf+1}$  invariant way ?**

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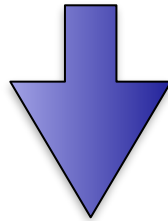
1. Coulomb moduli parameter is transformed.

2. Expanded in terms of instanton factor  $q = e^{-\frac{\beta}{2g^2}}$

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1. Use **invariant variable**  $\tilde{A} = q^{\frac{2}{8-N_f}} e^{-\beta a}$  **instead of**  $e^{-\beta a}$

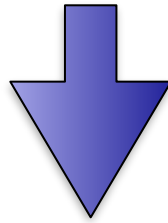
$$\left( \frac{1}{g^2} \rightarrow -\frac{1}{g^2}, \quad a \rightarrow a + \frac{1}{4g^2}, \quad \tilde{A} = q^{\frac{1}{4}} e^{\beta a} = e^{\beta(a + \frac{1}{8g^2})} \quad (N_f = 0) \right)$$



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2. Expand in terms of  $\tilde{A}$

# $E_{Nf+1}$ invariant Nekrasov partition function

$$Z(a, g, m_i; \epsilon_1, \epsilon_2) = Z_{pert}(a, m_i; \epsilon_1, \epsilon_2) \sum_{k=0}^{\infty} Z_k(a, m_i; \epsilon_1, \epsilon_2) q^k$$

**Original form**

'12 H-C Kim, S-S.Kim, K.Lee

'14 C.Hwang, J.Kim, S.Kim, J.Park

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$$= \sum_{n=0}^{\infty} \frac{\tilde{Z}_n(g, m_i; \epsilon_1, \epsilon_2)}{\mathbf{E}_{Nf+1} \text{ invariant}} \tilde{A}^k$$

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'14 C.Hwang, J.Kim, S.Kim, J.Park

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**New form**

'14 C.Hwang, J.Kim, S.Kim, J.Park

$$= \exp \left[ \sum_{k=0}^{\infty} \frac{1}{k} \sum_{n=1}^{\infty} \tilde{F}_n(kg, km_i; k\epsilon_1, k\epsilon_2) \tilde{A}^{nk} \right]$$

$$\equiv \text{PE} \left[ \sum_{n=1}^{\infty} \frac{\tilde{F}_n(g, m_i; \epsilon_1, \epsilon_2) \tilde{A}^n}{\mathbf{E}_{Nf+1} \text{ invariant}} \right]$$

# Nekrasov partition function for pure $SU(2)$

$$Z = \text{PE} \left[ \frac{\mathfrak{q} + \mathfrak{t}}{(1 - \mathfrak{q})(1 - \mathfrak{t})} \chi_2^{E_1} \tilde{A}^2 + \mathcal{O}(\tilde{A}^4) \right]$$

Character of  $E_1 = SU(2)$  :  $\chi_2^{E_1} = q^{\frac{1}{2}} + q^{-\frac{1}{2}}$ ,

$$\tilde{A} = q^{\frac{1}{4}} e^{\beta a}, \quad \mathfrak{q} = e^{-\beta \epsilon_1}, \quad \mathfrak{t} = e^{\beta \epsilon_2}$$

**Manifestly  $E_1$  invariant!!**

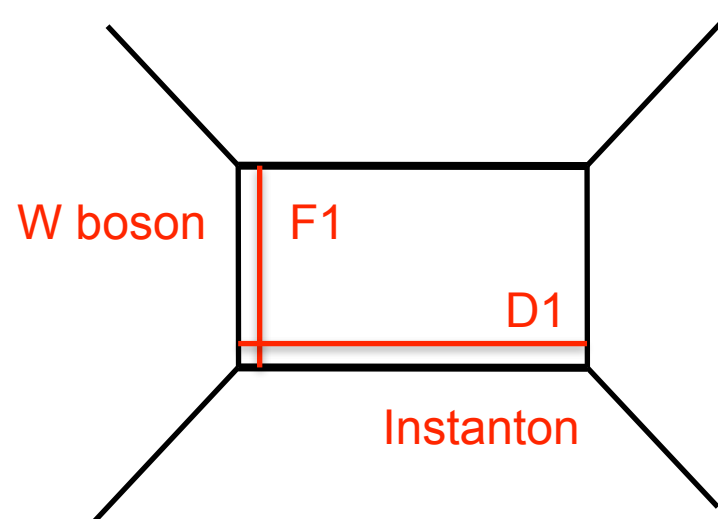
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# Nekrasov partition function for $N_f=1$

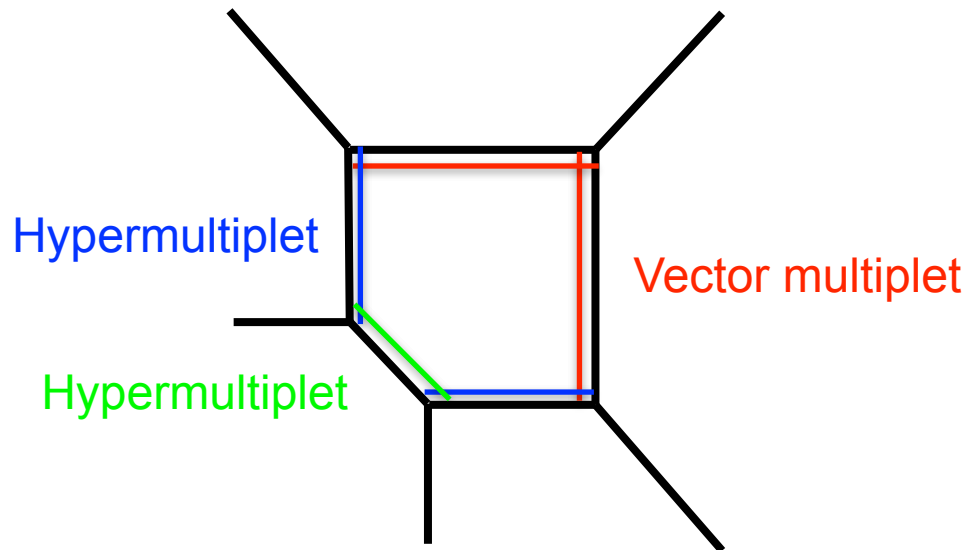
$$Z = \text{PE} \left[ -\frac{q^{\frac{1}{2}} t^{\frac{1}{2}}}{(1-q)(1-t)} \left( \chi_2^{SU(2)}(u_1) u_2^{-\frac{3}{7}} + u_2^{\frac{4}{7}} \right) \tilde{A} \right. \\ \left. + \frac{q+t}{(1-q)(1-t)} \chi_2^{SU(2)}(u_1) u_2^{\frac{1}{7}} \tilde{A}^2 + \mathcal{O}(\tilde{A}^3) \right]$$

$$E_2 = SU(2) \times U(1) \quad \begin{array}{ll} SU(2) : & u_1 = q^{\frac{1}{2}} e^{-\frac{1}{4}\beta m} \quad \chi_2(u_1) = u_1 + u_1^{-1} \\ U(1) : & u_2 = q^{-\frac{1}{2}} e^{-\frac{7}{4}\beta m} \quad \tilde{A} = q^{\frac{2}{7}} e^{-\beta a} \end{array}$$

# Nekrasov partition function for $N_f=1$

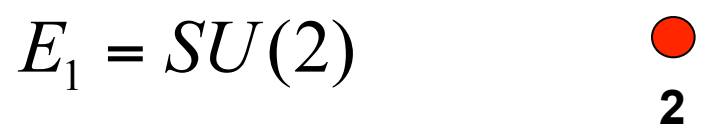
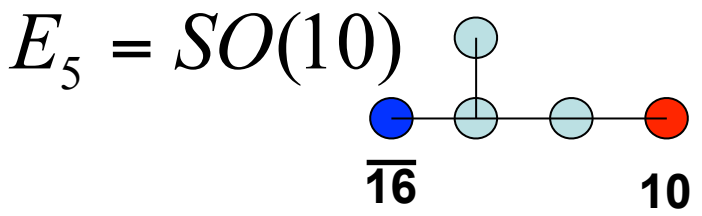
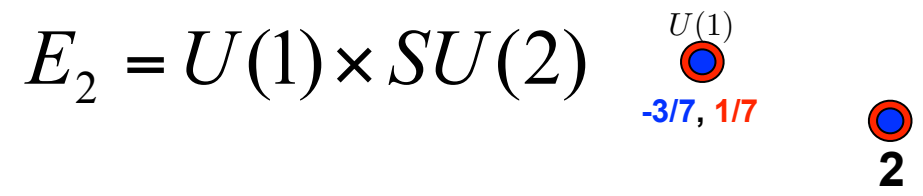
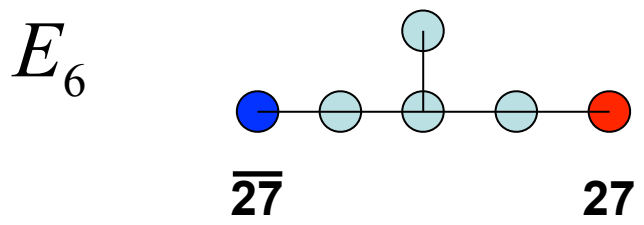
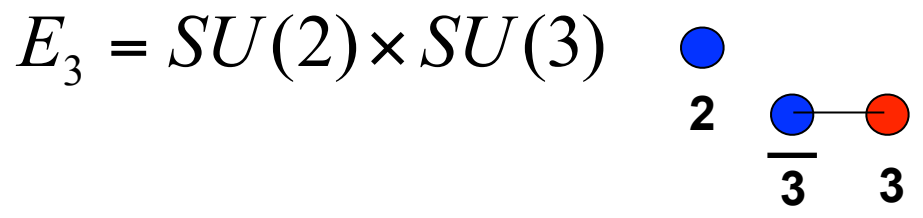
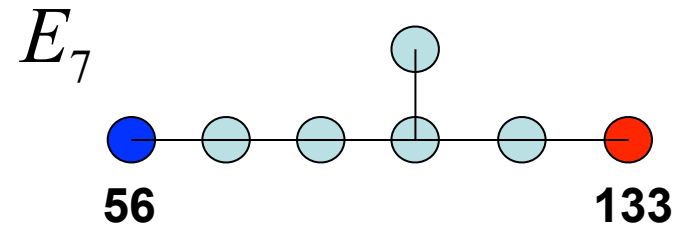
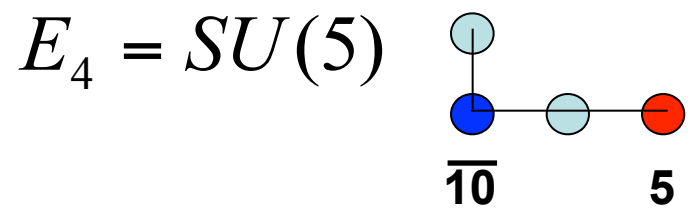
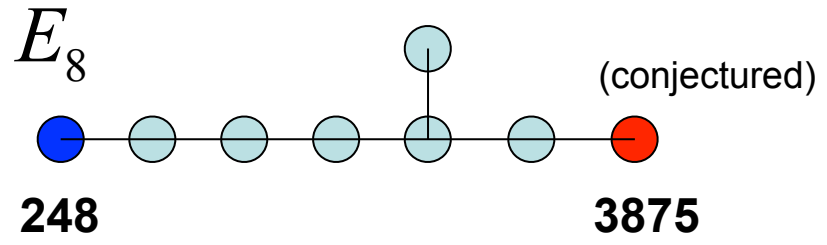
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The **vector multiplet** and the **hypermultiplet** are included in the fundamental representation of  $E_{N_f+1}$  corresponding to the following nodes



# Gopakumar-Vafa's expansion

$$Z = \text{PE} \left[ \sum_{C \in H_2(X, \mathbb{Z})} \sum_{j_L, j_R} \sum_{k_L = -j_L}^{j_L} \sum_{k_R = -j_R}^{j_R} \frac{M_C^{(j_L, j_R)} t^{k_L + k_R} q^{k_L - k_R}}{(t - t^{-1})(q - q^{-1})} Q^C \right]$$


$X$  : Calabi-Yau manifold

$Q^C = e^{-\int_C \omega}$ ,  $\omega$  : Kähler form

( $Q^C = e^{-2\beta a}, e^{-\beta(a-m)}, q^k e^{-2\beta a}, \dots$ )

$M_C^{(j_L, j_R)}$  : Refined Gopakumar-Vafa invariant

Gopakumar-Vafa '98  
Iqbal, Kozcaz, Vafa '07

**Nekrasov partition function**  **Set of integers**  $M_C^{(j_L, j_R)}$

**Non-Negative integer**

(After the convention change  $\tilde{A} \rightarrow -\tilde{A}$ )

**Consistent with the result from topological B-model**

[Huang, Klemm, Poretschkin '13]

# Summary

$$SO(2N_f) \times U(1) + \text{S-duality} = E_{N_f+1}$$

Nekrasov partition function is invariant

Refined Gopakumar-Vafa invariants from  
Nekrasov partition function agrees with  
topological B-model computation